



Simulations of flows over spill weirs with *Code_Saturne*

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- Project initiated by EDF R&D studying geometries of **spill weir**

Goal: enhance adherence, limit low pressures on the weir and delay dynamic flow separation up to higher upstream water heights.

- **Volume Of Fluid** method new in *Code_Saturne* V5.0.

- 1 VOF model
- 2 Thin weir
- 3 Creager weir
- 4 Piano Key (PK) Weir
- 5 References

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Starting with Navier-Stokes equation:

$$\begin{cases} \frac{\partial \rho}{\partial t} + \underline{\nabla} \cdot (\rho \underline{u}) = 0 \\ \frac{\partial \rho \underline{u}}{\partial t} + \underline{\nabla} \cdot (\rho \underline{u} \otimes \underline{u} - \underline{\underline{T}}) = \rho \underline{g} \end{cases} \quad (1)$$

with $\underline{\underline{T}}$ the stress tensor:

$$\underline{\underline{T}} = -(P + \frac{2}{3}\mu \underline{\nabla} \cdot \underline{u}) \underline{\underline{1}} + \mu (\underline{\underline{\nabla}} \underline{u} + (\underline{\underline{\nabla}} \underline{u})^T)$$

- An air volume fraction (often called void fraction) is defined: α .
- Mixing laws are used:

$$\rho = \alpha\rho_{air} + (1 - \alpha)\rho_{water} \quad (2)$$

$$\mu = \alpha\mu_{air} + (1 - \alpha)\mu_{water} \quad (3)$$

- By injecting (2) in the mass balance equation, we get:

$$\begin{aligned} \frac{\partial(\alpha\rho_{air} + (1 - \alpha)\rho_{water})}{\partial t} + \underline{\nabla} \cdot ([\alpha\rho_{air} + (1 - \alpha)\rho_{water}]\underline{u}) &= 0 \\ \Leftrightarrow (\rho_{air} - \rho_{water})\left[\frac{\partial\alpha}{\partial t} + \underline{\nabla} \cdot (\alpha\underline{u})\right] + \rho_{water}\underline{\nabla} \cdot (\underline{u}) &= 0 \end{aligned}$$

- and the following equivalence between the transport equation of the void fraction and the incompressibility equation:

$$\boxed{\frac{\partial\alpha}{\partial t} + \underline{\nabla} \cdot (\alpha\underline{u}) = 0 \Leftrightarrow \underline{\nabla} \cdot (\underline{u}) = 0} \quad (4)$$

VOF model:

$$\underline{\nabla} \cdot (\underline{u}) = 0$$

$$\frac{\partial \rho \underline{u}}{\partial t} + \underline{\nabla} \cdot (\rho \underline{u} \otimes \underline{u}) = \rho \underline{g} - \underline{\nabla} P + \underline{\nabla} \cdot (\mu (\underline{\underline{\nabla}} \underline{u} + (\underline{\underline{\nabla}} \underline{u})^T) - (\frac{2}{3} \mu \underline{\underline{\nabla}} \cdot \underline{u}) \underline{\underline{1}})$$

$$\frac{\partial \alpha}{\partial t} + \underline{\nabla} \cdot (\alpha \underline{u}) = 0$$

Momentum balance under its solved form, with $P^* = P - \rho_{air} \underline{g}$ the solved pressure:

$$\begin{aligned} \rho \frac{\partial \underline{u}}{\partial t} + \underline{\nabla} \cdot (\rho \underline{u} \otimes \underline{u}) - \underline{u} \underline{\nabla} \cdot (\rho \underline{u}) \\ = (\rho - \rho_{air}) \underline{g} - \underline{\nabla} P^* + \underline{\nabla} \cdot (\mu (\underline{\underline{\nabla}} \underline{u} + (\underline{\underline{\nabla}} \underline{u})^T) - (\frac{2}{3} \mu \underline{\underline{\nabla}} \cdot \underline{u}) \underline{\underline{1}}) \end{aligned}$$

VOF algorithm in *Code_Saturne*

- Prediction step: computation of $\tilde{\underline{u}}$

$$\rho \frac{\tilde{\underline{u}} - \underline{u}^n}{\Delta t} + \nabla \cdot ((\rho \underline{u})^n \otimes \tilde{\underline{u}}) - \tilde{\underline{u}} \nabla \cdot ((\rho \underline{u})^n) = (\rho - \rho_{air}) \underline{g} - \nabla P^n + \nabla \cdot (\mu (\underline{\nabla} \tilde{\underline{u}} + (\underline{\nabla} \tilde{\underline{u}})^T) - (\frac{2}{3} \mu \underline{\nabla} \cdot \tilde{\underline{u}}) \underline{\underline{1}})$$

- Correction step:

$$\delta P = P^{n+1} - P^n$$

and:

$$\begin{cases} \frac{\underline{u}^{n+1} - \tilde{\underline{u}}}{\Delta t} = -\frac{1}{\rho} \underline{\nabla}(\delta P) \\ \underline{\nabla} \cdot \underline{u}^{n+1} = 0 \end{cases}$$

yielding:

$$\underline{\nabla} \cdot (\tilde{\underline{u}}) = \underline{\nabla} \cdot \left(\frac{\Delta t}{\rho} \underline{\nabla}(\delta P) \right)$$

δP_{cell} then $(\underline{u}^{n+1} \cdot \underline{S})_{f_{ij}}$ are computed, and finally \underline{u}_{cell}^n is updated to $\underline{u}_{cell}^{n+1}$

- Resolution of the void fraction transport equation and update of $(\rho \underline{u} \cdot \underline{S})_f$ with an upwind scheme and of mixing properties ρ and μ

$$\frac{(\alpha_f^{n+1} - \alpha_f^n)}{\Delta t} + \underline{\nabla} \cdot (\alpha \underline{u})^{n+1} = 0$$

- Possibility to iterate over these steps to make $\tilde{\underline{u}}$ converge towards \underline{u}^{n+1} (in remaining terms appearing in the sum of prediction and correction equations).

Discretization over the cell Ω_i and between t_n and t_{n+1} :

$$\frac{\partial \alpha}{\partial t} \rightarrow \frac{|\Omega_i|}{\Delta t} (\alpha_i^{n+1} - \alpha_i^n)$$

$$\underline{\nabla} \cdot (\alpha \underline{u}) \rightarrow \sum_f [(1 - \theta)(\alpha_f \underline{S}_f \cdot \underline{u}_f)^n + \theta(\alpha_f \underline{S}_f \cdot \underline{u}_f)^{n+1}]$$

where:

f face of Ω_i

\underline{S}_f surface vector of face f

α_f face value of α (interpolated)

\underline{u}_f face value of \underline{u} (interpolated)

Discrete transport equation applying an implicit Euler time scheme ($\theta = 1$):

$$\frac{|\Omega_i|}{\Delta t} (\alpha_i^{n+1} - \alpha_i^n) + \sum_f (\alpha_f \underline{S}_f \cdot \underline{u}_f)^{n+1} = 0$$

Numerical scheme choice

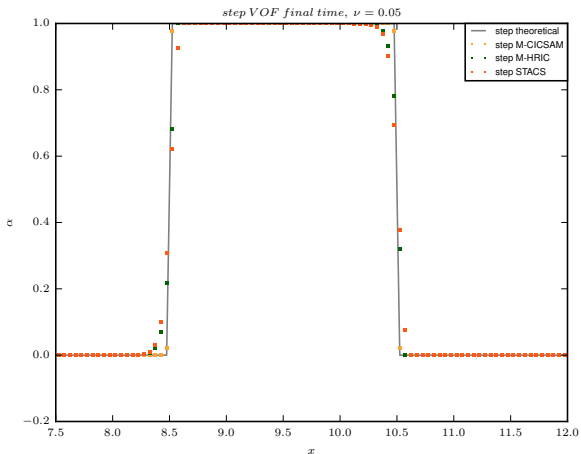
Necessary properties:

- α bounded between 0 and 1 - min/max principle.
- avoid interface diffusion.

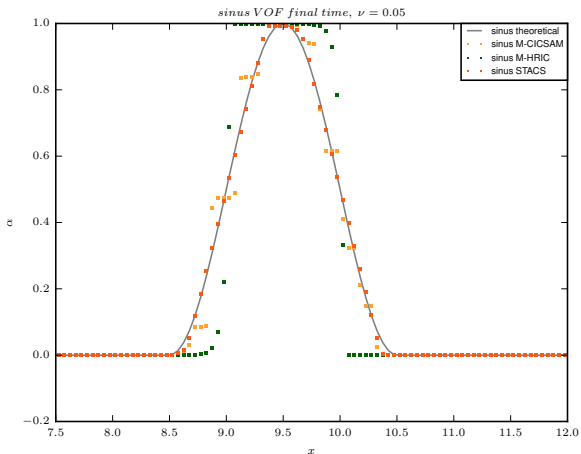
Numerical convection schemes implemented in *Code_Saturne* for VOF computation: STACS [2], M-HRIC [4], M-CICSAM [7]

Principle: blend a very compressive scheme and a Total Variation Diminishing scheme, with a weighting factor depending on angle between the free surface and the face normal, and on the Courant number Co .

Convection scheme comparisons - 1D convection of a step signal



Convection scheme comparisons - 1D convection of a sinus signal



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Simulation of a flow over a thin weir

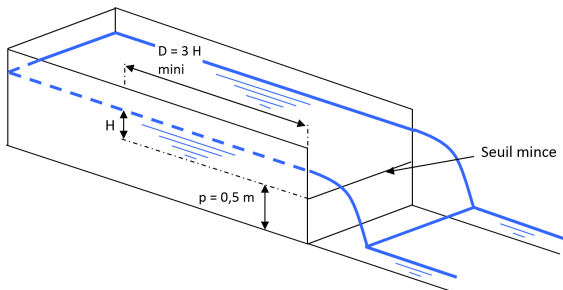


Figure: Sketch of a spilling way with a thin weir

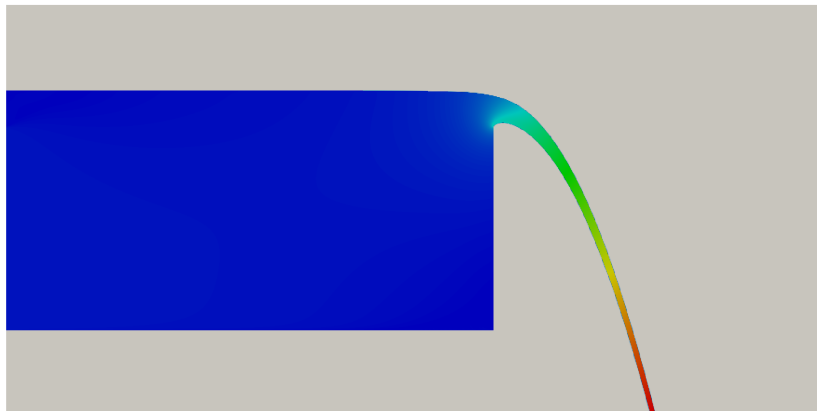


Figure: Velocity for $q = 0.1\text{m}^2/\text{s}$

Validation of *Code_Saturne* simulations by comparison with empirical laws:

- Height-Flow rate law:

$$Q = C_D b \sqrt{2g} H^{3/2}$$

$$C_D = 0.4023 \left(\frac{h_e}{H}\right)^{3/2} \left(1 + 0.135 \frac{h_e}{p}\right) \quad (\text{Rehbock [3]})$$

$$\text{with } h_e = H + 0.0011$$

- Profile of the lowest free surface of the water flow:

$$\frac{z}{H} = \frac{1}{2} \left(\frac{x}{H}\right)^{1.85} \quad (\text{Scimemi [5]})$$

$$\frac{z}{H} = 0.556 \left(\frac{x}{H}\right)^2 \quad (\text{De Marchi [5]})$$

Thin weir - comparison to Height-Flow rate laws

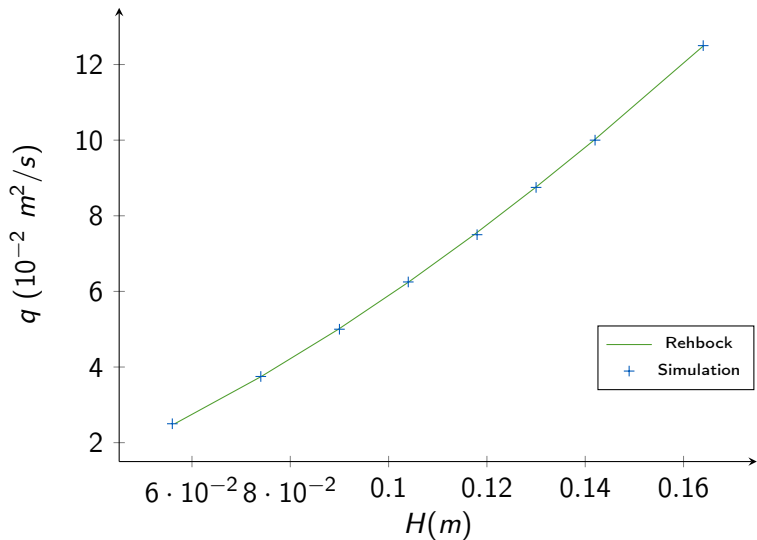
Free surface defined for the criterium $\alpha = 0.5$, for the 2nd mesh, the thickness of the free surface ($\alpha \in [0.001, 0.999]$) is not bigger than one cell height (2mm) for all flow rates.

Débit (m^2/s)	H_{th} (m)	H_{simu} (m) $\pm 2mm$	H_{simu} (m) $\pm 1mm$	Écart %	C_D Rehbock
0.025	0.0565	0.056	0.056	-0.85	0.4205
0.0375	0.0741	0.0726	0.074	-0.15	0.4196
0.05	0.0898	0.088	0.09	0.27	0.4198
0.0625	0.1041	0.104	0.104	-0.06	0.4203
0.075	0.1174	0.116	0.118	0.52	0.4210
0.0875	0.1299		0.13	0.06	0.4218
0.1	0.1418	0.14	0.142	0.13	0.4227
0.125	0.1641	0.162	0.164	-0.06	0.4245

Table: Values of the height H for several flow rates

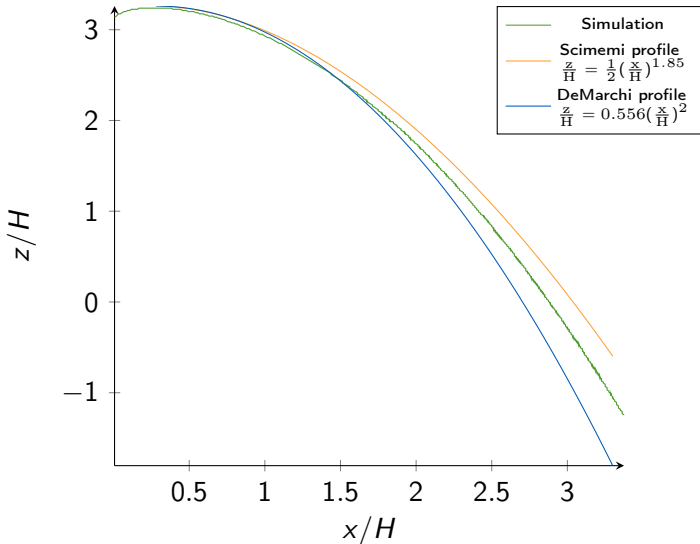
Thin weir - comparison to Height-Flow rate laws

Comparing to Rehbock law



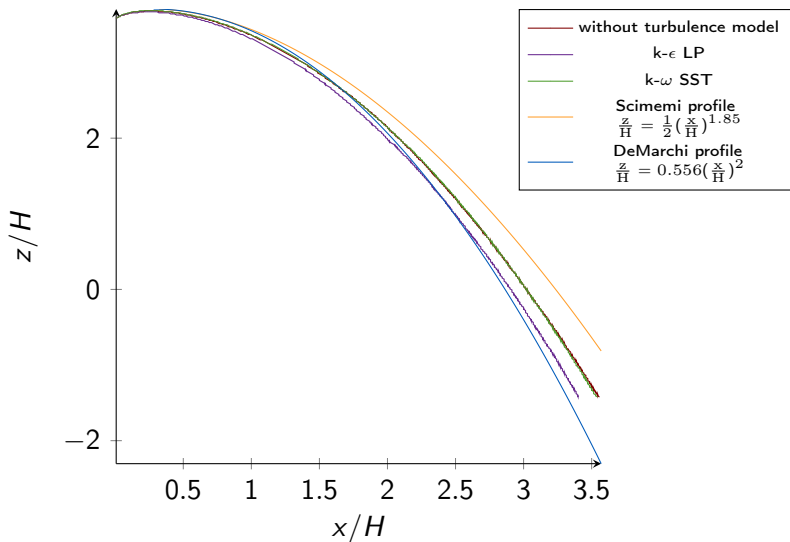
Thin weir - lowest free surface profile

Free surface for $q = 0.125 m^2 \cdot s^{-1}$



Thin weir - Turbulence model sensitivity

Free surface for $q = 0.1 \text{ m}^2 \cdot \text{s}^{-1}$



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Creager weir

Simulation of a flow spilling over a Creager weir
chosen (dimensioning) height: 20 cm (reduced size)

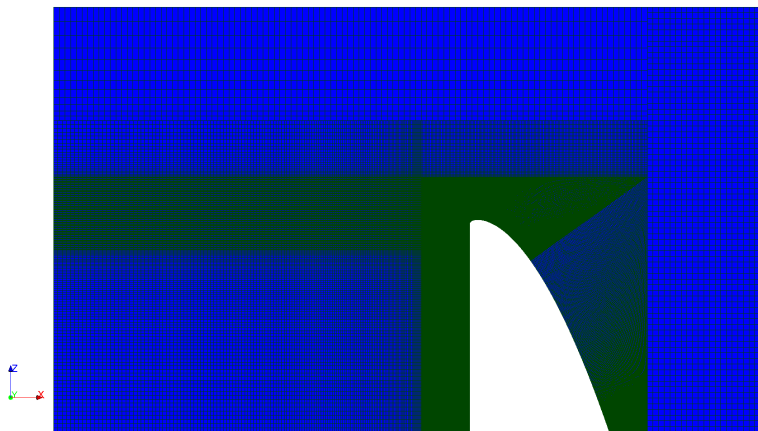


Figure: Mesh of the Creager weir

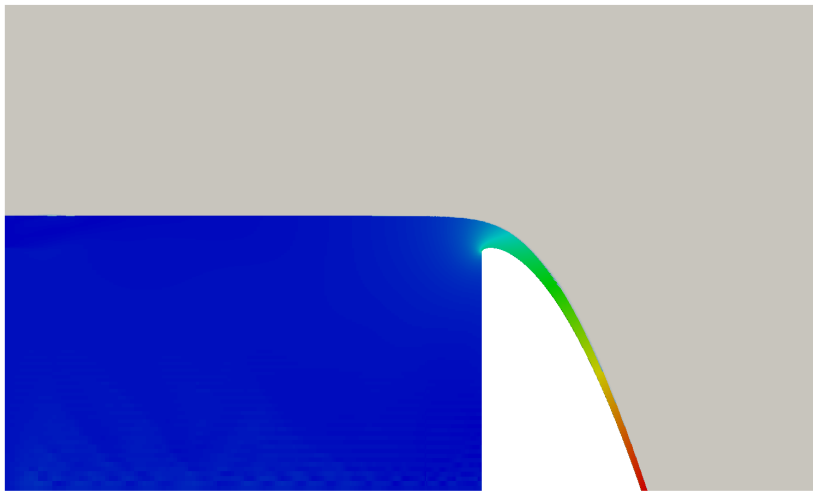


Figure: velocity field for $q = 0.2\text{m}^2/\text{s}$

Validation of the simulations by comparing with empirical laws:

- Height-flow rate:

$$Q = C_D b \sqrt{2g} H^{3/2}$$

$$C_D = 0.4950 \left(\frac{H}{H_D} \right)^{0.12} \quad \text{pour } (0.2 < \frac{H}{H_D} < 2) \quad (\text{Brudenell [6]})$$

- Free surface profile (Vischer and Hager [1]):

$$S = 0.75 \left(\chi^{1.1} - \frac{1}{6} X \right) \quad \text{for } (-2 < X/\chi^{1.1} < 2)$$

$$\text{with } S = s/H_D \quad \chi = H/H_D \quad X = x/H_D$$

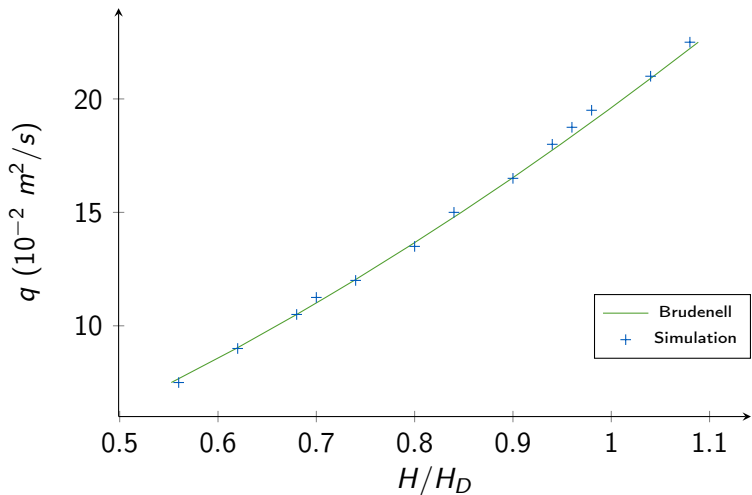
Creager weir - Height-Flow rate law

Flow rate (m^2/s)	H_{th} (m)	H_{simu} (m) $\pm 2mm$	Deviation %	C_D	H/H_D
0.075	0.1105	0.112	1.36	0.461	0.55
0.09	0.1237	0.124	0.27	0.467	0.62
0.105	0.1360	0.136	-0.01	0.473	0.68
0.1125	0.1419	0.14	-1.36	0.475	0.71
0.12	0.1477	0.148	0.21	0.477	0.74
0.135	0.1588	0.16	0.74	0.482	0.79
0.15	0.1695	0.168	-0.88	0.485	0.85
0.165	0.1798	0.18	0.13	0.489	0.90
0.18	0.1897	0.188	-0.89	0.492	0.95
0.1875	0.1945	0.192	-1.30	0.493	0.97
0.195	0.1993	0.196	-1.65	0.495	1.00
0.21	0.2086	0.208	-0.30	0.498	1.04
0.225	0.2177	0.216	-0.78	0.500	1.09
0.525	0.3673	0.36	-1.98	0.533	1.84
0.6	0.3989	0.39	-2.22	0.538	1.99

Table: Values of the height H - Surface: $\alpha = 0.5$ - $H_D = 0.2m$

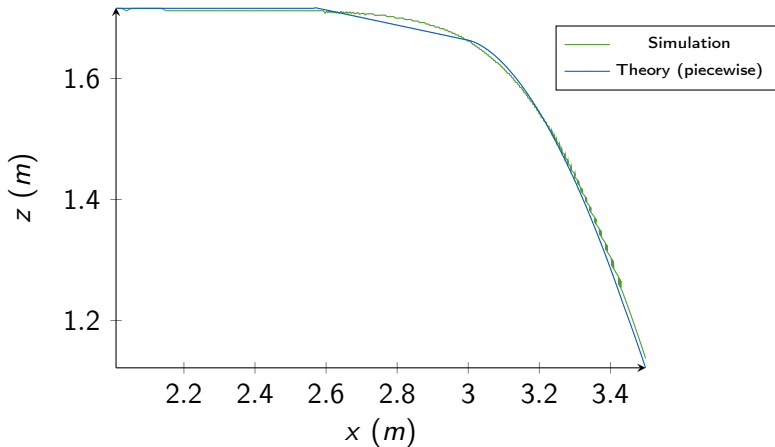
Creager weir - Height-Flow rate law

Comparing simulation results against Brudenell formula



Creager weir - free surface profile

Free surface for $q = 0.125 \text{ m}^2 \cdot \text{s}^{-1}$



Creager weir - large scale

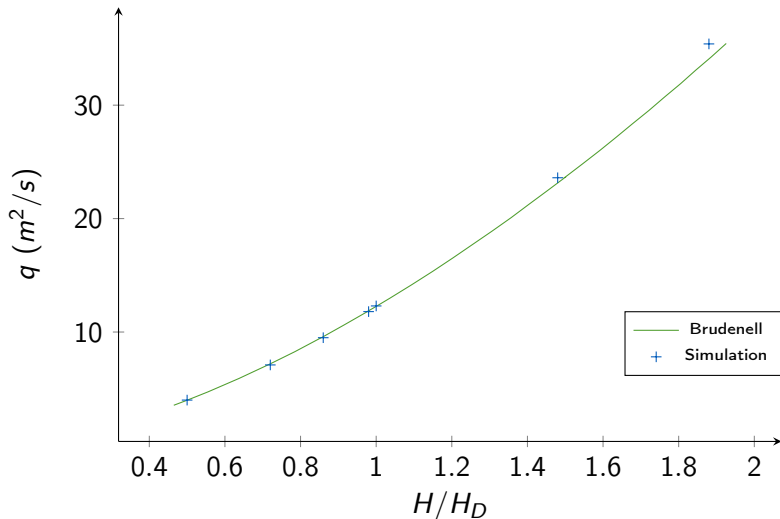
Real size: $H_D = 3.15m$ (for the comparison with *Flow3D* results on a validation case used by hydraulic engineers at CIH)

Uncertainty on height measurement: $\pm 0.063m$

Flow rate (m^2/s)	H_{th} (m)	H_{simu} (m)	Deviation %	H/H_D	$C_{D\ th}$	$C_{D\ simu}$	Deviation %
4.0	1.576	1.575	-0.1	0.50	0.4555	0.4555	-0.01
7.1	2.246	2.268	1.0	0.71	0.4753	0.4759	0.12
9.5	2.683	2.709	1.0	0.85	0.4856	0.4861	0.12
11.8	3.079	3.087	0.3	0.98	0.4936	0.4938	0.03
12.3	3.154	3.15	-0.1	1.00	0.4951	0.4950	-0.02
23.6	4.723	4.662	-1.3	1.50	0.5197	0.5188	-0.16
35.4	6.066	5.922	-2.4	1.93	0.5355	0.5340	-0.29

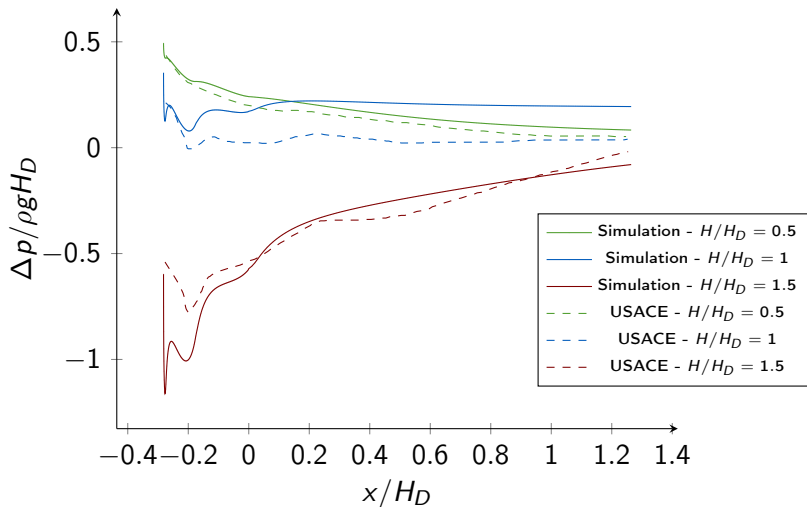
Table: Values of height H for several flow rates - $H_D = 3.15m$

Comparison of simulation with Brudenell formula



Large scale Creager - pressure along surface of the weir

Pressure in boundary cells



Large scale Creager - reduced Creager and flow separation

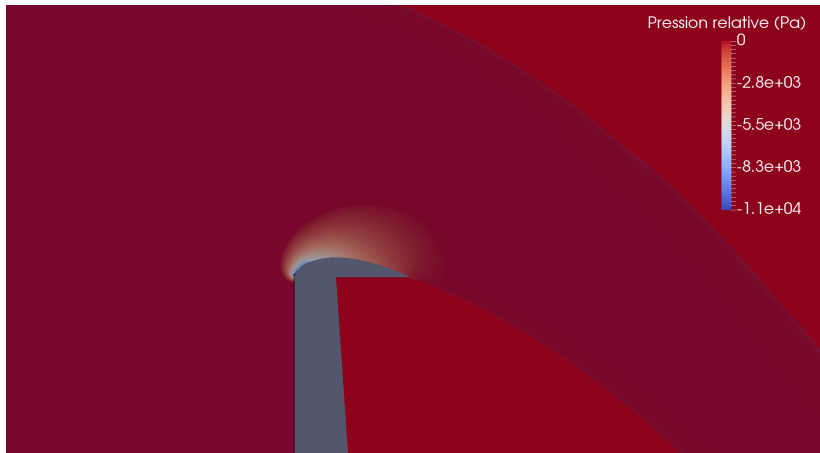


Figure: Pressure along reduced Creager weir for $H/H_D = 2.5$

Creager weir - 3D case

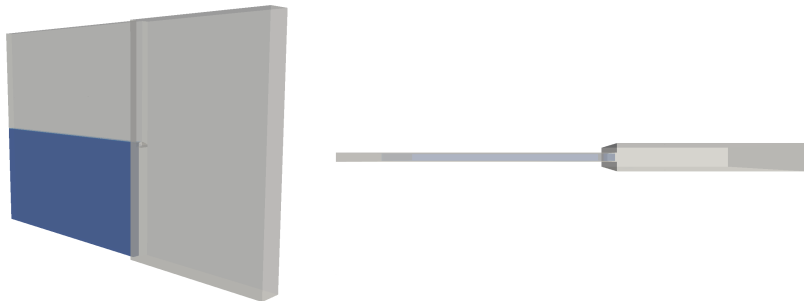


Figure: Geometry for the Creager weir 3D simulation - Initialisation

Creager weir - 3D case

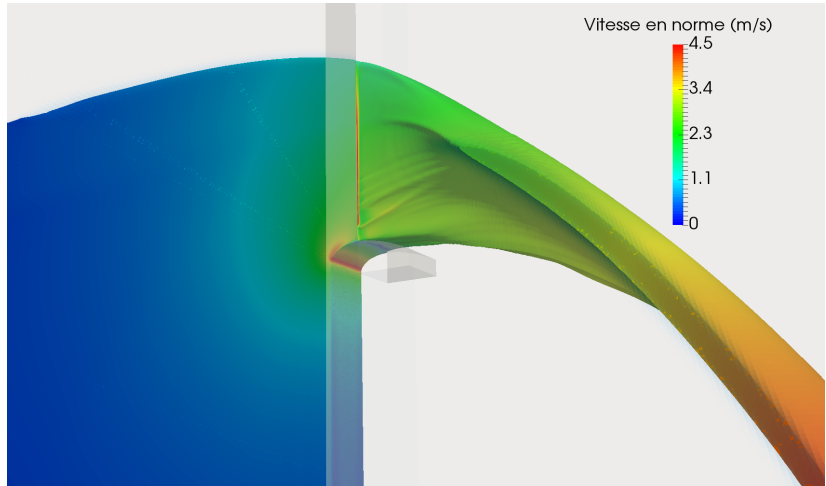


Figure: Flow separating from the weir for $H/H_D = 2.3$

Creager weir - 3D case

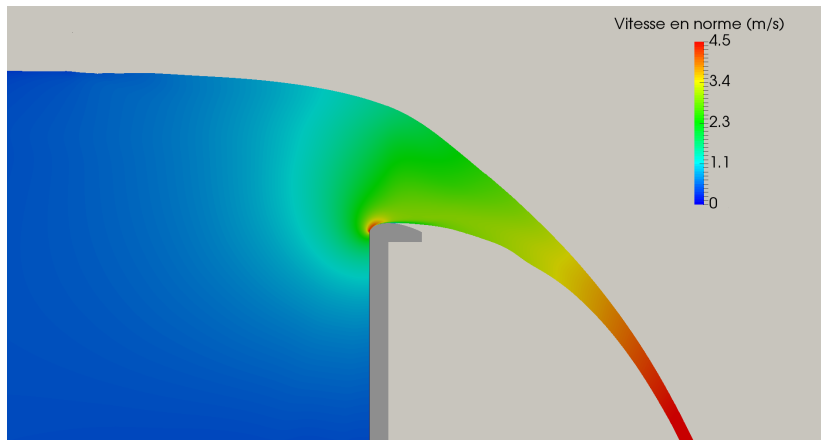


Figure: Slice view (middle plane along flow direction) - $H/H_D = 2.3$

Creager weir - 3D case

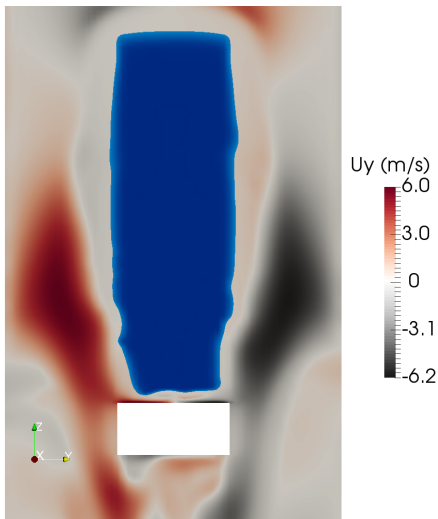


Figure: Slice with (y, z) plane at $x/H_D = 0.3$ - y velocity component - blue zone at the image center belongs to spilling flow - $H/H_D = 2.3$

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Piano Key (PK) Weir

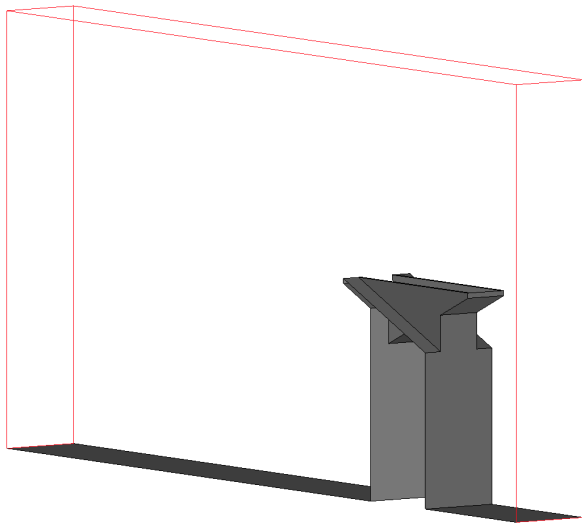


Figure: Geometry of a Piano Key Weir - height 4m



Figure: PK Weir - Mesh of the domain (min. cell size 5cm)

Simulation - Flow rate $15.6 \text{ m}^2/\text{s}$

PK Weir - comparison with height-flow rate law

Flow rate (m^2/s)	H_{exp} (m)	H_{simu} (m) $\pm 5cm$	Deviation %
3.5	0.5	0.55	10
8.2	1	1.2	20
15.6	2	2.3	15
22.4	3	3.3	10
25.5	3.5	3.9	11.4
35.5	5	5.1	2

Table: Comparison of the values of the height head H - experimental data : F. Lempérière

- Validation of the model for 2D overtopping flows.
- Low pressure area observed on Creager weirs and flow separation exhibited on 3D simulations of Creager weir.
- Improvements needed on PK weir computations (robustness issues - ongoing work).

- Explore new designs of spilling weirs.
- Optimize existing shapes of weir → ongoing internship, workflow in Salome.
- For the VOF module, other topic foreseen: waves propagation and overflow.

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References



Vischer D. and Hager W.
page 44, 1998.



M. Darwish and F. Moukalled.
Convective schemes for capturing interfaces of free-surface flows on unstructured grids.
Numerical Heat Transfer Fundamentals, 2006.



Carlier M.
1986.



Muzaferija S., Peric M., Sames P., and Schellin T.
A two-fluid navier-stokes solver to simulate water entry.
Proceedings of the Twenty-Second Symposium on Naval Hydrodynamics, 1999.



Sentürk.
1994.



Hager W. and Schleiss A.
15:173, 2009.



Di Zhang, Chunbo Jiang, Dongfang Liang, Zhengbing Chen, Yan Yang, and Ying Shi.
A refined vof algorithm for capturing sharp fluid interfaces on arbitrary meshes.
Journal of Computational Physics, 2014.